

NARESUAN UNIVERSITY
The Institute for Fundamental Study (IF)

**LINEAR ALGEBRA, LINEAR SPACES, EXAMPLES, SPECTRAL
DECOMPOSITION (M1)**

**FOURIER SERIES, FOURIER TRANSFORM, DIRAC DELTA,
DISTRIBUTIONS (M2)**

SUMMER SCHOOL
Einstein's Term 2024

Homework assignment
(Final update: May 20, 2024)

1st) Consider the following discrete orthonormal basis in $\mathcal{L}^2(\Omega)$, where Ω is the interval $[0, \ell]$:

$$\psi_a(x) = \sqrt{\frac{1}{\ell}} \exp\left(i \frac{2\pi a}{\ell} x\right), \quad a = \dots, -2, -1, 0, +1, +2, \dots \quad (1.1)$$

Every function $\psi(x) \in \mathcal{L}^2(\Omega)$ can be expanded in one and only one way in terms of the $\{\psi_a(x)\}$:

$$\psi(x) = \sum_{a=-\infty}^{\infty} C_a \psi_a(x). \quad (1.2)$$

As you know, the coefficients of the latter series are given by the formula $C_a = \langle \psi_a, \psi \rangle$. Also note that $\psi(x)$ satisfies the periodic boundary condition, i.e., $\psi(0) = \psi(\ell)$. (a) Try to obtain the typical Fourier series from the series given above, namely,

$$\psi(x) = A_0 + \sum_{a=1}^{\infty} A_a \cos\left(\frac{2\pi a}{\ell} x\right) + \sum_{a=1}^{\infty} B_a \sin\left(\frac{2\pi a}{\ell} x\right), \quad (1.3)$$

with the coefficients A_0 , A_a and B_a written in terms of the coefficients C_a , namely,

$$A_0 = \sqrt{\frac{1}{\ell}} C_0 = \frac{1}{\ell} \int_0^{\ell} dx \psi(x), \quad (1.4)$$

$$A_a = \sqrt{\frac{1}{\ell}} (C_a + C_{-a}) = \frac{2}{\ell} \int_0^{\ell} dx \cos\left(\frac{2\pi a}{\ell} x\right) \psi(x), \quad (1.5)$$

$$B_a = \sqrt{\frac{1}{\ell}} i(C_a - C_{-a}) = \frac{2}{\ell} \int_0^{\ell} dx \sin\left(\frac{2\pi a}{\ell} x\right) \psi(x). \quad (1.6)$$

(b) Demonstrate that $\psi(x)$ is real if and only if $C_{-a} = C_a^*$ (the asterisk * denotes the complex conjugate, as usual), and therefore, the coefficients A_0 , A_a and B_a are real.

2nd) Try to demonstrate the following integral representation for the Dirac delta:

$$\int_{\mathbb{R}} dx e^{ikx} = \int_{-\infty}^{+\infty} dx e^{ikx} = \lim_{a \rightarrow 0} \int_{-\infty}^{+\infty} dx e^{ikx} e^{-a|x|} = 2\pi \delta(k). \quad (2.1)$$

Note that, I am proposing you to add a regularizing factor to derive the expression. You will also need to use the following representation of the Dirac delta:

$$\lim_{a \rightarrow 0} \frac{1}{\pi} \frac{a}{a^2 + k^2} = \delta(k). \quad (2.2)$$

Note: In both limits, the parameter a approaches zero from the positive side.

3rd) As presented in class, if $\psi(x)$ is a (real or complex) function of the variable x , its Fourier transform $\phi(k)$, if it exist, is defined by

$$\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dx \psi(x) e^{-ikx} \equiv \text{FT}[\psi(x)], \quad (3.1)$$

and the inverse formula is:

$$\psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dk \phi(k) e^{+ikx} \equiv (\text{FT})^{-1}[\phi(k)]. \quad (3.2)$$

(a) Demonstrate the following property:

$$\text{FT}[\psi(cx)] = \frac{1}{|c|} \phi(k/c), \quad (3.3)$$

where c is a (real) constant. (b) In particular, $\text{FT}[\psi(-x)] = \phi(-k)$, and together with $\text{FT}[\psi(x)] = \phi(k)$ (see Eq. (3.1)), it follows that if the function $\psi(x)$ has a definite parity, its Fourier transform $\phi(k)$ has the same parity. Prove it!

4th) (a) Let us consider the following function:

$$\psi(x) = \frac{1}{\sqrt{2a}} [\Theta(x+a) - \Theta(x-a)], \quad (4.1)$$

where $a > 0$ and $\Theta(x)$ is the Heaviside step function, namely, $\Theta(x < 0) = 0$ and $\Theta(x > 0) = 1$. Find the Fourier transform of $\psi(x)$, namely,

$$\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} dx \psi(x) e^{-ikx}. \quad (4.2)$$

(b) Is the function $\psi(x)$ normalized to unity? In other words, is $\psi(x)$ a square-integrable function with norm equal to one? (c) But furthermore, is also $\phi(k)$ a square-integrable

function with norm equal to one? Hint: You may find the following integral useful:

$$\int_0^{+\infty} du \frac{\sin^2(u)}{u^2} = \frac{\pi}{2}. \quad (4.3)$$

5th) Prove the following result (in fact, it is a theorem): If $\phi(k)$ and $\varphi(k)$ are the respective Fourier transforms of the square-integrable functions $\psi(x)$ and $\chi(x)$, one has that

$$\int_{-\infty}^{+\infty} dx \chi^*(x) \psi(x) = \int_{-\infty}^{+\infty} dk \varphi^*(k) \phi(k). \quad (5.1)$$

A particular case of this result is the conservation of the norm

$$\int_{-\infty}^{+\infty} dx |\psi(x)|^2 = \int_{-\infty}^{+\infty} dk |\phi(k)|^2, \quad (5.2)$$

that is, a function and its Fourier transform have the same norm. This result is called the Parseval-Plancherel formula.

6th) (a) Demonstrate that the Fourier transform of the Heaviside step function $\Theta(x)$ is given by

$$\phi(k) = \text{FT}[\Theta(x)] = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} dx \Theta(x) e^{-ikx} = \frac{(-i)}{\sqrt{2\pi}} \lim_{a \rightarrow 0} \frac{1}{k - ia}, \quad (6.1)$$

(remember that $\Theta(x < 0) = 0$ and $\Theta(x > 0) = 1$). The limit can be evaluated using the following identity:

$$\begin{aligned} \lim_{a \rightarrow 0} \frac{1}{k \pm ia} &= \text{P.V.} \left(\frac{1}{k} \right) \mp i\pi \delta(k) \\ \left[\text{i.e., } \lim_{a \rightarrow 0} \int_{\mathbb{R}} dk \frac{f(k)}{k \pm ia} &= \text{P.V.} \int_{\mathbb{R}} dk \frac{f(k)}{k} \mp i\pi \int_{\mathbb{R}} dk f(k) \delta(k) \right. \\ &= \text{P.V.} \int_{\mathbb{R}} dk \frac{f(k)}{k} \mp i\pi f(0) \left. \right]. \end{aligned} \quad (6.2)$$

Certainly, $f(x)$ is a regular function at $k = 0$ (in these limits, the parameter approaches zero from the positive side), and P.V. represents the Cauchy principal value of the integral to the right of the symbol P.V., namely,

$$\text{P.V.} \int_{\mathbb{R}} dk \frac{f(k)}{k} \equiv \lim_{\epsilon \rightarrow 0} \int_{-\infty}^{-\epsilon} dk \frac{f(k)}{k} + \lim_{\epsilon \rightarrow 0} \int_{+\epsilon}^{+\infty} dk \frac{f(k)}{k} \quad (\epsilon > 0). \quad (6.3)$$

(b) Find the inverse Fourier transform of $\phi(k)$

$$(\text{FT})^{-1}[\phi(k)] = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} dk \phi(k) e^{+ikx}. \quad (6.4)$$

Did you find what you expected? Hint:

$$\frac{1}{\pi} \int_{\mathbb{R}} dk \frac{\sin(kx)}{k} = \text{sgn}(x) = 2\Theta(x) - 1. \quad (6.5)$$

7th) Let $\psi = \psi(x)$ be a function on which the operator \hat{p}^{-1} acts, where $\hat{p} = -i\hbar d/dx$ is the momentum operator. Let $\phi = \phi(x)$ be the result of this operation:

$$\phi(x) = \hat{p}^{-1} \psi(x). \quad (7.1)$$

(a) How does \hat{p}^{-1} act on the function $\psi(x)$? I will tell you how to answer this question. If $f(p)$ and $g(p)$ are, respectively, the Fourier transforms of $\psi(x)$ and $\phi(x)$, we can write the following relations:

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dp f(p) \exp\left(+i\frac{p}{\hbar}x\right) \Rightarrow f(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dx \psi(x) \exp\left(-i\frac{p}{\hbar}x\right) \quad (7.2)$$

and

$$\phi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dp g(p) \exp\left(+i\frac{p}{\hbar}x\right) \Rightarrow g(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dx \phi(x) \exp\left(-i\frac{p}{\hbar}x\right). \quad (7.3)$$

Show that

$$g(p) = p^{-1} f(p). \quad (7.4)$$

Multiplying the latter relation by $\exp(+ipx/\hbar)/\sqrt{2\pi\hbar}$, and then integrating over p , the following result is obtained:

$$\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dp g(p) \exp\left(+i\frac{p}{\hbar}x\right) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dp \frac{1}{p} f(p) \exp\left(+i\frac{p}{\hbar}x\right),$$

or even better

$$\phi(x) = \frac{i}{2\pi\hbar} \int_{-\infty}^{+\infty} dx' \left\{ \int_{-\infty}^{+\infty} dp \frac{1}{p} \sin\left[\frac{p}{\hbar}(x-x')\right] \right\} \psi(x') \quad (7.5)$$

(where expressions included in Eqs. (7.2) and (7.3) were used). Show that the formula (7.5) can be written as follows:

$$\phi(x) = \frac{i}{\hbar} \int_{-\infty}^x dx' \psi(x') - \frac{i}{2\hbar} \int_{-\infty}^{+\infty} dx' \psi(x'). \quad (7.6)$$

Note: Formula (7.4) shows that the function $g(p)$ has, in general, a pole at $p = 0$. To avoid this result, the condition $f(p = 0) = 0$ could be imposed, which implies that $\int_{-\infty}^{+\infty} dx' \psi(x') = 0$. Thus, the result in Eq. (7.6) can be written as follows:

$$\phi(x) = \frac{i}{\hbar} \int_{-\infty}^x dx' \psi(x'). \quad (7.7)$$

Clearly, the operator \hat{p}^{-1} is an integral operator of the form

$$\hat{p}^{-1} = \frac{i}{\hbar} \int_{-\infty}^x dx' (), \quad (7.8)$$

and acts on a function of x (which should be placed inside the parenthesis in (7.8) as a function of x'). (b) Final note: Verify, using formula (7.7), that $\hat{p} \phi(x) = \psi(x)$.

8th) In a given representation, or basis, a bounded operator can be represented by a matrix. If we change the representation, the operator will be represented by a different matrix. On the other hand, if we write the matrix in the basis of its own eigenvectors, then the matrix is diagonal. Write the Pauli matrix $\hat{X} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ in the basis of its eigenvectors. Which matrix did you find?

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